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SUFFIELD MEMORANDUM

NO. 1081 - Vol

A NONLINEAR SIX DEGREE-OF-FREEDOM

BALLISTIC AERIAL TARGET SIMULATION MODEL

VOLUME 1. THEORETICAL DEVELOPMENT (U)

by

A.B. Markov

PCN No. 21V10

February 1984

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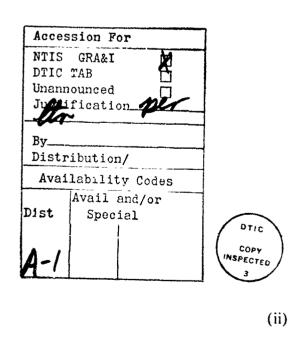
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A NONLINEAR SIX DEGREE-OF-FREEDOM BALLISTIC AERIAL TARGET SIMULATION MODEL (U)

VOLUME 1. THEORETICAL DEVELOPMENT
VOLUME 2. COMPUTER PROGRAM USER MANUAL

by

A.B. Markov

ABSTRACT

Six degree-of-freedom, rigid body equations of motion are described suitable for modeling the dynamic characteristics of multistaged, free-flight, ballistic rockets such as the DRES developed aerial targets CRV7/BATS and ROBOT-9. These equations of motion form the core of a FORTRAN simulation software package called BALSIM. This package allows for modeling of vehicle thrust and structural asymmetries, time-varying mass and inertia characteristics, variable wind conditions, nonstandard atmospheric conditions, stage failures, and different rocket motor types. The BALSIM package has been written in IBM FORTRAN IV and has been tested on the IBM 3033 computer with the H-extended compiler. It is currently being adapted for use with the VAX11/780 and Honeywell DPS-8/70C computers.

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Figure A2.1 Definition of Reference Frame F_{TH_1}

LIST OF SYMBOLS

The following is a list of important symbols. Some symbols, which are defined in the text and are used only once, or are secondary quantities related to a primary quantity that is apparent from the text or from the notation conventions to follow are not included. Numbers in parentheses refer to equations.

- Accodynamic force vector applied to the vehicle not including thrust forces
- a Vehicle centre-of-mass acceleration vector relative to inertial space
- a Speed of sound
- b Reference length (fuselage diameter for ballistic rocket vehicles)
- C_D (Total drag)/(½ ϱ V²S)
- $C_{L_{fin}}$ (Fin lift normal to fin chord)/($\frac{1}{2}\varrho V^2S$)

$$(C_{L_{\alpha}})_{i,i}$$
 $\frac{\partial C_{L_{fin}}}{\partial \alpha} \mid e$

$$C_i$$
 $L_{A_B}/(\frac{1}{2}\varrho V^2Sb)$

$$\begin{array}{c|c} C_{I_p} & \frac{2V}{b} \frac{\partial C_I}{\partial p_B} e \end{array}$$

$$C_{i_{\vec{0}\vec{k}}}$$
. $\frac{\partial C_i}{\partial \delta_{fin}} = e$

$$C_m = M_{A_B}/(\frac{1}{2} \varrho V^2 St)$$

$$C_{m_q}$$
 $\frac{2V}{b} \frac{\partial C_m}{\partial q_B} = e$

$$C_n = N_{A_B}/(1/2 \varrho V^2 Sb)$$

$$C_{n_r}$$
 $\frac{2V}{b} \frac{\partial C_n}{\partial r_B} = 2$

 $C_y = Y_{A_R}/(\frac{1}{2}\varrho V^2 S)$

 $C_{y_{fit}}$ (Aerodynanic force along y-axis of F_B generated by

pseudo fin)/(½ QV2S)

 $C_{y_{ofin}}$ $C_{y_{fin}}$ for $\alpha_{fin} = 0$

 $(C_{y\alpha})_{fin}$ $\frac{\partial C_{y_{fin}}}{\partial \alpha_{fin}} = e$

 $C^{B}_{\nu\beta}$ $\frac{\partial C^{B}_{\nu}}{\partial \beta}$ e

 $C_z^B = Z_{A_B}/(1/2 \varrho V^2 S)$

 $C_{z_{t,n}}^{B}$ (Aerodynamic force along z-axis of F_{B} generated

by pseudo fin)/($\frac{1}{2}\varrho V^2S$)

 $C_{z_{ofin}}$ $C_{z_{fin}}$ for $\alpha_{fin} = 0$

 $C_{z\alpha}$ $\frac{\partial C_z}{\partial \alpha}$ e

 $(C_{z_{\alpha}})_{fin}$ $\frac{\partial C_{z_{fin}}}{\partial \alpha_{fin}}$ e

D Total drag

 e_{fin} The body-fin interference factor of the pseudo fin

External force vector acting on the vehicle centre-of-mass

 F_B Body-fixed reference frame, see Figure 2

F_B' Modified body-fixed reference frame, see Figure 2

F₁ Inertial reference frame, see Figure 1

F_L Launcher inertial reference frame, see Figure 1

F_R Structural body-fixed reference frame, see Figure 2

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F_T	Stand-off inertial reference frame, see Figure !
g	Acceleration due to gravity
go	Nominal sea level acceleration due to gravity
<u>h</u>	Vehicle angular momentum vector about centre-of-mass
h _{ASL}	Altitude of vehicle centre-of-mass above sea level
I_{sp_i}	Specific impulse of i-th rocket motor
I_{xx}^B , I_{yy}^B ,	Vehicle moments of inertia about its centre-of-mass written as components in F_B
$(L_{A_B}, M_{A_B}, N_{A_B})$	Aerodynamic moment components in F_B not including thrust moments, about centre-of-mass
$(L_{\tau_B},M_{\tau_B},N_{\tau_B})$	Thrust moment components in F ₃ about centre-of-mass
$(L_{T_R}, M_{T_B}, N_{T_B})_{cg}$	See equation (2.6.1)
$(L_{T_B}, M_{T_B}, N_{T_B})_{nz}$	See equation (2.6.1)
M_A	Aerodynamic moment vector acting about the vehicle centre-of- mass not including thrust moments
M_T	Thrust moment vector acting about the vehicle centre-of-mass
m	Vehicle total mass
m _{em}	Airframe mass
$(\mathbf{m}_{Me})_{\iota}$	Mass of i-th rocket motor less propellant
M_{PL}	Payload mass
$(m_{PR})_{i}$	Mass of i-th rocket motor's propellant
N _M	1 otal number of motors
p_A	Atmospheric pressure
(p_B, q_B, r_B)	Angular velocity components of F_B with respect to F_I written as components in F_B

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q_{D}	Dynamic presure, ½ ρV ²
<u>R</u>	Position vector of vehicle centre-of-mass relative to F ₁
$\underset{\tau}{\mathbb{R}}$	Position vector of vehicle centre-of-mass relative to F_T
$\overset{\tau}{\underset{\tau'}{\Rightarrow}}$	Position vector of R_T relative to F_I
R	See equation (2.9,5)
Γ_E	Radius of the Earth to the nominal sea level datum plane
S	Reference area (fuselage cross-sectional area for ballistic rocket vehicles)
S	Distance the vehicle has moved along the launcher from its initial position [see equation (2.7,3)]
S _G	Launch rail guide length (see Figure 5)
Ţ	Thrust vector
T_A	Atmospheric temperature
T.	Thrust of the i-th rocket motor
(u_B, v_B, w_B)	Components of V in F_B
$(\mathbf{u}_{B_E}, \mathbf{v}_{B_E}, \mathbf{w}_{B_E})$	Components of V_E in F_B
$(\mathbf{u}_{B_g}, \mathbf{v}_{B_g}, \mathbf{w}_{B_g})$	Components of $\underset{\longrightarrow}{W}$ in F_B
V	Airspeed vector
$\overset{\mathbf{V}}{\longrightarrow}_{\mathcal{E}}$	Velocity vector of the vehicle centre-of-mass with respect to F,
V	Magnitude of V
V_{xz}	$\sqrt{u_B^2 + w_B^2}$
W	Wind velocity vector with respect to F,
(W_1, W_2, W_3)	Components of W in F_t

W_{fin}	Pseudo fin airspeed component normal to its chord plane, see equation (2.4,11)
$(X_{A_B}, Y_{A_B}, Z_{A_B})$	Aerodynamic force components in F_B not including thrust contribution acting at the vehicle centre-of-mass
$(X_{T_B}, Y_{T_B}, Z_{T_B})$	Thrust vector components in F_B
(X_{ac}, Y_{ac}, Z_{ac})	Coordinates of the vehicle aerodynamic centre in F_R
$(X_{ac_{fin}}, y_{ac_{fin}}, Z_{ac_{fin}})$	Coordinates of the aerodynamic centre of the pseudo fin in F_R
(x_{cg}, y_{cg}, z_{cg})	Coordinates of the vehicle centre-of-mass in F_R
(X_{em}, Y_{em}, Z_{em})	Coordinates of the airframe (empty) centre-of-mass in F_R
(x_T, y_Y, z_T)	Components of R_T in F_T [see equation (2.10,3)]
(x_I, y_I, z_I)	Components of R in F_I
$[(\mathbf{X}_{Me})_{\iota}, (\mathbf{y}_{Me})_{\iota}, (\mathbf{Z}_{Me})_{\iota}]$	Coordinates of the centre-of-mass of the empty motor case of the i-th rocket motor
(X_{PL}, Y_{PL}, Z_{PL})	Coordinates of the payload centre-of-mass in F_R
$[(X_{PR})_{i}, (Y_{PR})_{i}, (Z_{PR})_{i}]$	Coordinates of the centre-of-mass of the propellant of the i-th rocket motor
α	Angle of attack of vehicle [see equation (2.4,6a)]
$lpha_{fin}$	Angle of attack of pseudo fin's chord plane [see equation (2.4,12b)]
β	Sideslip angle of vehicle [see equation (2.4,6b)]
δ_{fin}	Cant angle of pseudo fin (see Figure 4)
$ heta_{\scriptscriptstyle B}$	Elevation Euler angle of F_B
ξA	Aspect elevation angle of vehicle relative to F_T
$\xi_{\it E}$	Aspect azimuth angle of vehicle relative to F_T

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Air density Q θ_B Euler bank angle for F_B $\theta_{\scriptscriptstyle fin}$ Cylindrical coordinate of pseudo fin (see Figure 4) Euler azimuth angle for F_B ψ_B Angular velocity vector of F_B with respect to F_I ω_B NOTATION CONVENTIONS F_A Reference frame A X Vector quantity X $X \times Y$ Vector cross product of X and Y A matrix X X X^T The transpose of X or the components of a vector X expressed in the reference frame F_T (con will determine which interpretation is intended) X expressed as components in F_A X_A X'A $F_{B'}$ variable that is analogous to the variable X in F_{B} A column matrix x X

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quasisteady equilibrium condition

X,

A quantity x whose value is computed for an aerodynamic

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VOLUME 1. THEORETICAL DEVELOPMENT (U)

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1. INTRODUCTION

In 1979 DRES began development work towards modifying the U.S. army BATS target for use with the higher specific impulse CRV7 rocket motors under the auspices of a joint US-Canadian TTCP agreement. This original work has lead to a number of DRES initiated activities including the modification of the target to permit multistaging and the development of an all CRV7, multistaged vehicle referred to as ROBOT-9 (Figure 1). As well, support equipment has been developed for target operation in moderately heavy seas, i.e. ir conditions typical of Canadian coastal waters (the ROBOT System development history is summarized in more detail in Reference 1).

To support the development of these free-flight targets, computer simulation programs were required that predicted the dynamic characteristics of the vehicles. Of particular importance were accurate predictions of basic performance parameters (e.g. range and flight-time), wind effects, effect of nonstandard atmospheric conditions and the dynamic effects of launching from a moving ship on a finite, nonzero length launcher.

No existing software was available to DRES that performed all of these tasks while conveniently permitting some configuration variation. As a result, in the period September, 1980 to December, 1981, a six degree-of-freedom simulation package was written, debugged, and tested at DRES. This package was applied to the evaluation of the CRV7/BATS and ROBOT-9 performance characteristics and safety-envelopes. It was coded in IBM FORTRAN and has been used on the IBM 3033 computer with the H-extended compiler. The package is currently being installed on a VAX 11/780 computer for use with a FORTRAN 77 compiler, and will be adapted for use with the Honeywell DPS-8/70C computer. The package has been designated BALSIM.

It is the intent of this report to provide documentation of BALSIM in sufficient detail to permit users familiar with FORTRAN to run the program. Chapter 2 develops the dynamic model and summarizes its limitations. Chapter 3 describes the BALSIM package in general terms. Finally, Volume 2 is intended as an essentially self-contained userbook for the package, and includes a listing of all program modules.

2. DYNAMIC MODEL

This section summarizes the key features of the dynamic model programmed into the BALSIM package. The basic six degree-of-freedom equations are derived in the following sections.

2.1 Fundamental Assumptions

Several overall simplifying assumptions have been made in the derivation of the equations of motion. They are valid for ballistic rocket vehicles that have rigid structures and relatively short ranges, i.e. less than 100 km (50 nm).

The assumptions are as follows:

- 1. The Earth is flat and any Earth-fixed reference frame is inertial.
- 2. The vehicle is a rigid body.
- 3. There are no control surfaces.

Assumption 3 may be readily relaxed by adding the appropriate control terms into the equations of motion.

2.2 Reference Frames, Rotation Matrices and Angular Velocities

In the general case both an Earth-fixed (say F_E) and an inertial reference frame (say F_I) must be defined. Because of the first simplifying assumption of the previous section, these reference frames become identical. Only F_I will be used here. Thus let F_I be an Earth-fixed inertial reference frame whose origin is at the launch site, whose x-axis points along the projection of the nominal launch trajectory onto the Earth's surface and whose z-axis is nominally downwards (see Fig. 1). The y-axis follows from the right hand rule.

A second, inertial Earth-fixed reference frame that is useful is the laurcher reference frame F_L . The origin of F_L is located at the launch site with the x-axis pointing in the launch direction and the z-axis being nominally downwards (see Fig. 1). The y-axis follows from the right hand rule.

A third Earth-fixed reference frame which is occasionally required is that of a reference frame F_T placed at some distance from the launch site. This reference frame may be used to compute the rocket aspect angle presentations (e.g. from the training ship). Since the location and orientation of this reference will depend on the particular application, it is defined only generally in Figure 1.

Since the aerodynamic forces are most conveniently expressed with respect to the vehicle, a body-fixed reference frame F_B will also be used. The origin of F_B is located at the vehicle centre-of-mass. In vehicles that are axisymmetric, the x-axis is on the axis of symmetry and points forward through the nose. Otherwise the x-axis points in the nominal faunch direction. The z-axis is nominally downward, while the y-axis follows from the right hand rule. This reference frame and some associated aerodynamic angles are shown in Figure 2.

For cases where the rocket vehicle mass characteristics are axisymmetric, the aerodynamic forces are independent of the vehicle's roll attitude, and the thrust forces are axisymmetric, the body-fixed reference frame need not spin with the vehicle. Thus a reference frame F_B is defined which is identical to F_B except that it does not rotate with the vehicle about the axis of symmetry. Initially F_B and F_B will coincide.

A third body-fixed reference frame that is useful in specifying the vehicle's mass, inertia and configuration characteristics is a reference frame F_R whose origin is located on a nose datum plane on the vehicle. If the vehicle is axisymmetric, then the origin of F_R

is on the axis of symmetry, and its x-axis points towards the rear of the vehicle on the axis of symmetry. Otherwise, the origin of F_R may be any convenient location on the nose datum plane, and the x-axis nominally points towards the rear of the vehicle. The z-axis is nominally upward and the y-axis follows from the right hand rule (see Figure 2).

In the following, use is made of a number of notation conventions. In particular if L_{BA} denotes a rotation matrix relating the components of a vector \underline{V} expressed in $F_A(\underline{V}_A)$ to the components of the same vector in $F_B(\underline{V}_B)$, then

$$\underline{\mathbf{v}}^{B} = \underline{\mathbf{L}}_{BA} \underline{\mathbf{v}}^{A} \tag{2.2,1}$$

The following definitions, geometric relationships, and matrices, will be employed in the presentation of the equations of motion.

1. The rotation matrix relating F_I and F_B :

$$\underline{\mathbf{L}}_{BI} = \begin{bmatrix} \cos\theta_B \cos\psi_B & \cos\theta_B \sin\psi_B & -\sin\theta_B \\ \sin\phi_B \sin\theta_B \cos\psi_B & \sin\phi_B \sin\psi_B \\ -\cos\phi_B \sin\psi_B & +\cos\phi_B \cos\psi_B \\ \cos\phi_B \sin\theta_B \cos\psi_B & \cos\phi_B \sin\theta_B \sin\psi_B \\ +\sin\phi_B \sin\psi_B & -\sin\phi_B \cos\psi_B \end{bmatrix}$$

$$= \begin{bmatrix} 1_{BI_{11}} \end{bmatrix}$$

$$(2.2,2a)$$

$$= \begin{bmatrix} 1_{BI_{12}} \end{bmatrix}$$

$$(2.2,2b)$$

where ϕ_B , θ_B , and ψ_B are the Euler angles defined by Etkin (Reference 3).

The rotation matrix relating F_I and $F_{B'}$ follows from (2.2,2a) by substituting $\psi_{B'}$, $\theta_{B'}$ and $\phi_{B'}$ for ψ_{B} , θ_{B} and ϕ_{B} respectively, and will be denoted $L_{B'I}$.

2. The rotation matrix relating F_L and F_I :

$$\underline{\mathbf{L}}_{LI} = \begin{bmatrix} \cos\theta_{B_o} & 0 & -\sin\theta_{B_o} \\ 0 & 1 & 0 \\ \sin\theta_{B_o} & 0 & \cos\theta_{B_o} \end{bmatrix}$$

$$= [\mathbf{l}_{LI_{i,j}}]$$

$$(2.2,3b)$$

3. The rotation matrix relating F_T and F_I :

$$\underline{L}_{TI} = \begin{bmatrix} \cos \psi_T & \sin \psi_T & 0 \\ -\sin \psi_T & \cos \psi_T & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1_{TI_{II}} \end{bmatrix}$$

$$(2.2,4b)$$

4. The angular velocity of the vehicle with respect to F_I written as components in F_B and F_{B}' :

$$\underline{\omega}^{B} = (p_{B}, q_{B}, r_{B})^{T}$$
 (2.2,5a)

$$\underline{\omega}^{B'} = (0, q_{B'}, r_{K'})^T$$
 (2.2,5b)

5. The angular rate cross-product matrices (Reference 3) in F_B and $F_{B'}$:

$$\frac{\widetilde{\omega}^{B}}{\omega} = \begin{bmatrix} 0 & -\mathbf{r}_{B} & \mathbf{q}_{B} \\ \mathbf{r}_{B} & 0 & -\mathbf{p}_{B} \\ -\mathbf{q}_{U} & \mathbf{p}_{B} & 0 \end{bmatrix}$$
 (2.2,6a)

$$\frac{\widetilde{\omega}^{B}}{\widetilde{\omega}^{B}} = \begin{bmatrix}
0 & -r_{B} & q_{B} \\
r_{B} & 0 & -p_{B} \\
-q_{U} & p_{B} & 0
\end{bmatrix}$$

$$\frac{\widetilde{\omega}^{B}}{\widetilde{\omega}^{B}} = \begin{bmatrix}
0 & -r_{B'} & q_{B'} \\
r_{B'} & 0 & 0 \\
-q_{B'} & 0 & 0
\end{bmatrix}$$
(2.2,6a)

6. The airspeed vector of the vehicle written as components in F_B and $F_{B'}$:

$$\underline{V}^{3} = (u_{B}, v_{B}, w_{B})^{T}$$
 (2.2,7a)

$$\underline{V}^{B'} = (u_B, v_{B'}, w_{B'})^T \tag{2.2.7b}$$

7. The groundspeed vector of the vehicle with respect to F_I written as components in F_B , $F_{B'}$ and F_I :

$$\underline{\mathbf{V}}_{E}^{B} = (\mathbf{u}_{B_{E}}, \mathbf{v}_{B_{E}}, \mathbf{w}_{B_{E}})^{T}$$
 (2.2,8a)

$$\underline{V}_{E}^{B'} = (\mathbf{u}_{B_{E}}, \mathbf{v}_{B_{E}}', \mathbf{w}_{B_{E}}')^{T}$$
 (2.2,8b)

$$\underline{V}_{E}^{I} = (\dot{x}_{I}, \dot{y}_{I}, \dot{z}_{I})^{T}$$
 (2.2,8c)

8. The aerodynamic angles (see Figure 2):

$$\alpha = \arctan(w_B/u_B) \tag{2.2,9}$$

$$\beta = \arctan (v_B/V_{xx}) \tag{2.2,10}$$

where

$$V_{xx} = \sqrt{u_B^2 + w_B^2}$$
 (2.2,11)

$$V = \sqrt{u_B^2 + v_B^2 + w_B^2}$$
 (2.2,12)

Here α is the angle of attack of the x-axis of F_B , β is the sideslip angle of the x-axis of F_B , V is the airspeed, and V_{xz} is the magnitude of the airspeed vector component along the x-z plane of F_B .

Analogous angles to α and β may be written in terms of $F_{B'}$ components by direct substitution of $F_{B'}$ quantities for F_{B} quantities.

9. The geometric relationships (see Figure 2):

$$u_B = V \cos \beta \cos \alpha \tag{2.2,13a}$$

$$v_B = V \sin \beta \tag{2.2,13b}$$

$$W_B = V \cos \beta \sin \alpha \tag{2.2,13c}$$

10. The wind velocity with respect to F_I written as components in F_I , F_B , and $F_{B'}$:

$$W^{I} = (W_1, W_2, W_3)^{T}$$
 (2.2,14a)

$$W_{B}^{B} = (u_{B_{g}}, v_{B_{g}}, w_{B_{g}})^{T}$$
 (2.2,14b)

$$\underline{\mathbf{W}}^{B'} = (\mathbf{u}_{B_g}, \mathbf{v}_{B_g}', \mathbf{w}_{B_g}')^T \tag{2.2,14c}$$

The acceleration due to gravity written as components in F₁:

$$g^{T} = (0, 0, g)^{T}$$
 (2.2,15)

12. The aerodynamic forces (not including thrust forces) written as components in F_B and F_{B}' :

$$\underline{\mathbf{A}}_{A}^{B} = (\mathbf{X}_{A_{B}}, \mathbf{Y}_{A_{B}}, \mathbf{Z}_{A_{B}})^{T}$$
 (2.2,16a)

$$A_A^{B'} = (X_{A_B}, Y_{A_B'}, Y_{A_B'})^T$$
 (2.2.16b)

13. The aerodynamic moments (not including thrust moments) written as components in F_B and $F_{B'}$:

$$\underline{M}_{A}^{B} = (L_{A_{B}}, M_{A_{B}}, N_{A_{B}})^{T}$$
 (2.2,17a)

$$M_A^{B'} = (L_{A_B}, M_{A_B'}, M_{A_B'})^T$$
 (2.2,17b)

14. The inertia matrix of the vehicle with respect to its centre-of-mass expressed in F_B (see Etkin, Reference 3) and $F_{B'}$:

$$\underline{\mathbf{I}}^{B} = \begin{bmatrix}
I_{xx}^{B} & -I_{xy}^{B} & -I_{xz}^{B} \\
-I_{xy}^{B} & I_{yy}^{B} & -I_{yz}^{B} \\
-I_{xx}^{B} & -I_{yz}^{B} & I_{zz}^{B}
\end{bmatrix} (2.2,18a)$$

$$\underline{\mathbf{I}}^{B'} = \begin{bmatrix}
I_{xx}^{B} & 0 & 0 \\
0 & I_{yy}^{B'} & 0 \\
0 & 0 & I_{yy}^{B'}
\end{bmatrix} (2.2,18b)$$

$$\underline{\mathbf{I}}^{B'} = \begin{bmatrix}
I_{xx}^{B} & 0 & 0 \\
0 & I_{yy}^{B'} & 0 \\
0 & 0 & I_{yy}^{B'}
\end{bmatrix} (2.2,18b)$$

15. The total thrust forces written as components in F_B and $F_{B'}$:

$$\underline{\mathbf{T}}^{B} = (\mathbf{X}_{T_{B}}, \mathbf{Y}_{T_{B}}, \mathbf{Z}_{T_{B}})^{T}$$
 (2.2,19a)

$$\underline{\mathbf{T}}^{B'} = (\mathbf{X}_{T_B}, \mathbf{Y}_{T_B'}, \mathbf{Y}_{T_B'})^T$$
 (2.2,19b)

16. The total thrust moments as components in F_B and $F_{B'}$:

$$\mathbf{M}_{T}^{B} = (\mathbf{L}_{T_{B}}, \mathbf{M}_{T_{A}}, \mathbf{N}_{T_{B}})^{T}$$
 (2.2,20a)

$$M_T^{B'} = (L_{T_B}, M_{T_B'}, M_{T_B'})^T$$
 (2.2,20b)

2.3 Newton-Euler Development of the General Equations of Motion

Newton-Euler techniques begin with the fundamental equations (Reference 3)

$$\mathbf{F} = \mathbf{ma} \tag{2.3,1}$$

and

$$\underline{\mathbf{h}} = \underline{\mathbf{M}} \tag{2.3,2}$$

a is the acceleration vector of the body centre of mass relative to an inertial reference frame, h is the angular momentum of the body about its centre-of-mass, F is the external force vector acting at the centre-of-mass and M is the external moment vector about the centre-of-mass. F may be written

$$\underline{\mathbf{F}} = \mathbf{m} \dot{\underline{\mathbf{V}}}_{E} \tag{2.3,3}$$

where V_E is the velocity vector of the vehicle with respect to F_I and m is mass of the vehicle. An expression for h follows from the fundamental relationship

or

where \underline{r} is the position vector of an element of mass dm of the body with respect to its centre-of-mass (see Figure 3), $\underline{\omega}_B$ is the angular velocity vector of F_B with respect to F_I , '.' when applied to a vector represents rate of change with respect to F_B and 'o' when applied to a vector represents rate of change with respect to F_B (see Reference 4 for a more thorough discuss on of vector differentiation). Equation (2.3,4b) may be written in matrix notation as (replacing $\underline{\omega}_B \times \underline{r}$ by $-\underline{r} \times \underline{\omega}_B$ and dropping the subscript 'B' on $\underline{\omega}_B$ for the sake of brevity)

^{*} Superscripts on matrix quantities refer to the reference frame in which the components of the matrix are expressed. Overscore '\circ' refers to the matrix equivalent of the vector cross-product.

$$\underline{\mathbf{h}}_{B} = \int_{\text{mass}} \left[\frac{\overset{\sim}{\mathbf{r}}^{B}}{\overset{\sim}{\mathbf{r}}^{B}} \cdot \frac{\overset{\sim}{\mathbf{r}}^{B}}{\overset{\sim}{\mathbf{r}}^{B}} \frac{\overset{\sim}{\mathbf{r}}^{B}}{\overset{\sim}{\mathbf{r}}^{B}} \underline{\omega}^{B} \right] d\mathbf{m}$$
 (2.3,5)

But

$$\underline{\dot{\mathbf{f}}}^{B} = \underline{\mathbf{0}} \tag{2.3.6}$$

for a rigid body. Since $\underline{\omega}^B$ is a constant with respect to the integration in (2.3,5), it follows that

$$h^B = I^B \omega^B \tag{2.3.7}$$

where

$$\underline{\mathbf{I}}^{B} = -\int_{mass} \underline{\underline{\mathbf{r}}}^{B} \underline{\underline{\mathbf{r}}}^{B} d\mathbf{m}$$
 (2.3,8)

 \underline{I}^{B} is, by convention, given by (2.2,18a).

The externally applied force \underline{F} is made up of an aerodynamic component \underline{A} , a thrust component \underline{T} and a gravitational component mg such that

$$\underline{F} = \underline{A} + \underline{T} + \underline{mg} \tag{2.3,9}$$

Substituting (2.3,9) into (2.3,1), the vector force equation becomes

$$m \underbrace{V_F} = \underbrace{A} + \underbrace{T} + mg \tag{2.3,10}$$

The externally applied moment $\underline{\underline{M}}$ is made up of an aerodynamic component $\underline{\underline{M}}_A$ and a thrust component $\underline{\underline{M}}_r$ such that

$$\underline{\mathbf{M}} = \underline{\mathbf{M}}_A + \underline{\mathbf{M}}_T \tag{2.3,11}$$

Substituting (2.3,11) into (2.3,2) yields

$$\dot{\underline{\mathbf{h}}} = \underline{\mathbf{M}}_A + \underline{\mathbf{M}}_T \tag{2.3,12}$$

Other than the gravitational force, the dominant forces and moments acting on the aircraft are the to aerodynamic causes and are largely determined by its orientation and configuration. It is accordingly advantageous to write the matrix equations of motion with respect to a body-fixed reference frame. This reference frame is chosen to be F_B . Furthermore, this choice does not introduce any gravitational moments since the origin of F_B and the centre-of-mass of the vehicle coincide.

Thus the matrix force and moment equations become

$$m(\underbrace{\mathbf{V}_{E}^{B}} + \underbrace{\mathbf{\omega}^{B}}_{\mathbf{\Sigma}} \underbrace{\mathbf{V}_{E}^{B}}) = \underbrace{\mathbf{A}^{B}} + \underbrace{\mathbf{T}^{B}}_{\mathbf{E}} + m \underbrace{\mathbf{L}_{BI}}_{\mathbf{E}} \mathbf{g}^{I} \qquad (2.3,13)$$

and

$$\dot{\underline{\mathbf{h}}}^{B} + \overset{\sim}{\omega}^{B} \dot{\underline{\mathbf{h}}}^{B} = \underline{\mathbf{M}}^{B}_{A} + \underline{\mathbf{M}}^{B}_{T} \tag{2.3,14}$$

Substituting for \underline{h}_B from (2.3,7), the moment equation becomes

$$\underline{\mathbf{I}}^{B} \dot{\underline{\omega}}^{B} + \dot{\underline{\mathbf{I}}}^{B} \underline{\omega}^{S} + \overset{\sim}{\underline{\omega}}^{B} \underline{\mathbf{I}}^{B} \underline{\omega}^{B} = \underline{\mathbf{M}}_{A}^{B} + \underline{\mathbf{M}}_{T}^{B}$$
 (2.3,15)

Writing out the equation (2.3,13) in scalar form yields

$$m(\dot{u}_{B_E} + q_B w_{B_E} - r_B v_{B_E}) = X_{A_B} + X_{T_B} + mgl_{Bl_{13}}$$
 (2.3,16a)

$$m(\dot{v}_{3_E} + r_B u_{B_L} - p_B w_{B_E}) = Y_{A_B} + Y_{T_B} + mgl_{BI_{23}}$$
 (2.3,16b)

$$m(\dot{w}_{B_E} + p_B I_B - q_B u_{B_E}) = Z_{A_B} + Z_{T_B} + mgl_{BI_{33}}$$
 (2.3,16c)

for the force equations, and

$$I_{xx}^{B}\dot{p}_{B} - I_{xy}^{B}\dot{q}_{B} - I_{xx}^{B}\dot{r}_{E} + I_{xx}^{B}p_{L} - I_{xy}^{B}q_{B} - I_{xx}^{B}r_{B} + I_{yz}(r_{B}^{2} - q_{B}^{2}) + (I_{zz}^{B} - I_{yy}^{B})r_{B}q_{B} + I_{yy}^{B}r_{B}p_{B} - I_{xz}^{B}q_{B}p_{B} = L_{A_{B}} + L_{T_{B}}$$

$$(3.2,17a)$$

$$-I_{xy}^{B}\dot{p}_{B} + I_{yy}^{B}\dot{q}_{B} - I_{yz}^{B}\dot{r}_{B} - \dot{I}_{xy}^{B}p_{B} + \dot{I}_{yy}^{B}q_{B} - \dot{I}_{yz}^{B}r_{F} + I_{xz}^{B}(p_{B}^{2} - r_{B}^{2}) + (I_{xx}^{2} - I_{zz}^{B})p_{B}r_{B} + I_{yz}^{B}p_{B}q_{B} - I_{xy}^{B}r_{B}q_{B} = M_{A_{B}} + M_{T_{B}}$$

$$(3.2,17b)$$

$$-I_{xz}^{B}\dot{p}_{B}-I_{yz}^{B}\dot{q}_{B}+I_{zz}^{B}\dot{r}_{B}-\dot{I}_{xz}^{B}p_{B}-\dot{I}_{yz}^{B}q_{B}+\dot{I}_{zz}^{B}r_{B}+I_{xy}^{B}(q_{B}^{2}-p_{B}^{2}) +(I_{yy}^{B}-I_{xx}^{B})p_{B}q_{B}+I_{xz}^{B}q_{B}r_{B}-I_{yz}^{B}p_{B}r_{B}=N_{A_{B}}+N_{T_{B}}$$
(3.2,17c)

for the moment equations.

Kinematic equations are also required for the linear and rotational position of the aircraft. The linear position equations follow from

$$\underline{V}_{E}^{I} = \underline{L}_{IB} \underline{V}_{E}^{B} \tag{2.3,18}$$

The rotational position equations are the Euler angle rate equations and are derived in Etkin (Reference 3). The resulting scalar kinematic equations of motion are thus seen to be

$$\dot{\mathbf{x}}_{I} = \mathbf{1}_{BI_{11}} \mathbf{u}_{B_{E}} + \mathbf{1}_{BI_{21}} \mathbf{v}_{B_{E}} + \mathbf{1}_{BI_{31}} \mathbf{w}_{B_{E}} \tag{3.2,19a}$$

$$\dot{y}_{I} = l_{BI_{12}} u_{B_{E}} + l_{BI_{22}} v_{3_{E}} + l_{BI_{32}} w_{B_{E}}$$
 (3.2,19b)

$$\dot{z}_{I} = l_{BI_{13}} u_{B_{E}} + l_{BI_{23}} v_{B_{E}} + l_{BI_{33}} w_{B_{E}}$$
 (3.2,19c)

for linear position, and

$$\dot{\phi}_B = p_B + q_B \sin \phi_B \tan \theta_B + r_B \cos \phi_B \tan \theta_B \qquad (2.3,20a)$$

$$\dot{\theta}_B = q_B \cos \phi_B - r_B \sin \phi_B \qquad (2.3,20b)$$

$$\dot{\psi}_B = [q_B \sin \phi_B + r_B \cos \phi_B] \sec \theta_B \qquad (2.3,20c)$$

for angular position.

It should be stressed that the variables $(u_{B_E}, v_{B_E}, w_{B_E})$ in equations (2.3,16), (2.3,17) and (2.3,19) are the body-axes components of the vehicle's ground velocity vector. This is not the same as the equations developed in Reference 2 where airspeed vector components in body-axes are used. Also, no assumptions have been made, up to this point, regarding vehicle planes of symmetry and the symmetry of the thrust and aerodynamic forces and moments. Finally, it should be noted that no assumptions have been made about the mass and inertia characteristics of the vehicle, i.e. in general

$$\dot{m} \neq 0$$

and

$$\dot{\mathbf{I}}_{ij}^{B} \neq 0$$

This, too, is an added feature not present in the equations developed in Reference 2.

A number of simplifications may be added to these equations if certain symmetry conditions are satisfied. If the vehicle has mass symmetry about the xy and xz planes, then

$$I_{xz}^B = I_{xy}^B = I_{yz}^B = 0 (2.3,21)$$

If the vehicle mass characteristics are also axisymmetric, then in addition to (2.3,21) we also have

$$I_{yy}^{B} = I_{zz}^{B} (2.3,22)$$

Applying assumption (2.3,21) to the moment equations (2.3,17a) through to (2.3,17c) results in the simplified set of equations

$$\dot{p}_B = \left[-\dot{I}_{xx}^B p_B - (I_{zz}^B - I_{yy}^B) r_B q_B + L_{A_B} + L_{T_B} \right] / I_{xx}^B$$
(2.3,23a)

$$\dot{q}_B = \left[-\dot{I}_{yy}^B q_B - (I_{xx}^B - I_{zz}^B) p_B r_B + M_{A_B} + M_{T_B} \right] / I_{yy}^B$$
 (2.3,23b)

$$\dot{r}_B = [-\dot{I}_{zz}^B r_B - (I_{yy}^B - I_{zz}^B) p_B q_B + N_{A_B} + N_{T_B}] / I_{zz}^B$$
 (2.3,23c)

Applying the axisymmetry assumption (2.3,22) to those equations simplifies (2.3,23a) even further by eliminating the $(I_{zz}^B - I_{yy}^B)r_Bq_B$ term, i.e.

$$\dot{p}_B = \left[-\dot{I}_{xx}^B p_B + L_{A_B} + L_{T_B} \right] / I_{xx}^B$$
(2.3,24)

If we now make the assumptions that the thrust forces are axisymmetric, that the vehicle mass characteristics are axisymmetric, and that the vehicle aerodynamic characteristics are independent of the vehicle's roll orientation, then we may take advantage of the simplifications that will result to the equations of motion by expressing them in the reference frame $F_{B'}$ rather than in F_{B} . Recall that in Section 2.2 we defined $F_{B'}$ as being identical to F_{B} except that it does not rotate with the vehicle about the axis of symmetry. As a result, if

$$\underline{\omega}_{B}' = p_{B}' \underline{i}_{B} + q_{B}' \underline{j}_{B} + r_{B}' \underline{k}_{B} \qquad (2.3.25)$$

is the angular velocity vector of \mathbf{F}_{B} relative to the inertial reference frame \mathbf{F}_{I} expressed

as components in F_{B} , then the latter condition implies that

$$p_{B'} = 0 (2.3,26)$$

for all t.

Also in general we will have

$$q_B' \neq q_B \tag{2.3,27a}$$

$$\mathbf{r}_{B}' \neq \mathbf{r}_{B} \tag{2.3,27b}$$

For the purpose of generating the aerodynamic forces, however, q_B' and r_B' may be treated as being interchangeable with q_B and r_B respectively because of the assumption that the aerodynamic forces are independent of the roll orientation. However, the resulting aerodynamic forces are now expressed in F_B' axes rather than in F_B axes.

There are also aerodynamic forces that are generated due to the vehicle's rolling angular velocity p_B . The latter is not identical to p_B' , and thus an equation of motion for p_B will still have to be retained.

Formally the equations of motion in F_B' may be obtained from the equations of motion in F_B by setting $p_B = 0$ in all equations except the p_B equation, replacing $(u_{B_E}, v_{B_E}, w_{B_E})$ with $(u_{B_E}, v_{B_E}, w_{B_E})$, (q_B, r_B) with (q_B', r_B') , $(\phi_B, \theta_B, \psi_B)$ with $(\phi_B', \theta_B', \psi_B')$, (M_{A_B}, N_{A_B}) with (M_{A_B}', M_{A_B}') , (M_{T_B}, N_{T_B}) with (M_{T_B}', M_{T_B}') , $(X_{A_B}, Y_{A_B}, Y_{A_B}, Z_{A_B})$ with $(X_{A_B}, Y_{A_B}', Y_{A_B}')$, $(X_{T_B}, Y_{T_B}, Z_{T_B})$ with $(X_{T_B}, Y_{T_B}', Y_{T_B}')$, and applying the assumptions discussed previously. Finally, the p_B equation from the F_B equations is retained.

The resulting equations of motion are as follows:

$$m(\dot{u}_{B_E} + q_B' w_{B_E}' - r_B' v_{B_E}') = X_{A_B} + X_{T_B} + mgl_{BI_1}'$$
 (2.3,28a)

$$m(\dot{v}'_{B_E} + r'_B u_{B_E}) = Y'_{A_B} + Y'_{T_B} + mgl'_{BI_{23}}$$
 (2.3,28b)

$$m(\dot{w}'_{B_F} - q_B u_{B_F}) = Y'_{A_B} + Y'_{T_B} + mgl'_{Bl_{22}}$$
 (2.3,28c)

$$\dot{p}_B = \left[-\dot{I}_{xx}^B p_B + L_{AB} + L_{TB} \right] / I_{xx}^B \tag{2.3,29a}$$

$$\dot{q}_B' = [-\dot{I}_{yy}^B q_B' + M_{A_B}' + M_{T_B}']/I_{yy}^B$$
 (2.3,29b)

$$\dot{\mathbf{r}}_{B}' = \left[-\dot{\mathbf{I}}_{vv}^{B} \mathbf{r}_{B}' + \mathbf{M}_{AB}' + \mathbf{M}_{TB}' \right] / \mathbf{I}_{vv}^{B}$$
 (2.3,29c)

$$\dot{\mathbf{x}}_{I} = \mathbf{1}'_{BI_{11}} \mathbf{u}_{B_{E}} + \mathbf{1}'_{BI_{21}} \mathbf{v}'_{B_{E}} + \mathbf{1}'_{BI_{31}} \mathbf{w}'_{B_{E}} \tag{3.2,30a}$$

$$\dot{y}_{I} = l'_{BI_{12}} u_{B_E} + l'_{BI_{22}} v'_{B_E} + l'_{BI_{32}} w'_{B_E}$$
 (3.2,30b)

$$\dot{z}_{I} = 1'_{BI_{13}} u_{B_{E}} + 1'_{BI_{23}} v'_{B_{E}} + 1'_{BI_{33}} w'_{B_{E}}$$
 (3.2,30c)

$$\frac{\dot{\phi}_B'}{\phi_B'} = q_B' \sin \phi_B' \tan \theta_B' + r_B' \cos \phi_B' \tan \theta_B' \qquad (2.3,31a)$$

$$\theta_B' = q_B' \cos \phi_B' - r_B' \sin \phi_B' \qquad (2.3,31b)$$

$$\psi_B' = [q_B' \sin \phi_B' + r_B' \cos \phi_B'] \sec \theta_B \qquad (2.3,31c)$$

2.4 Aerodynamic Model

The equations of motion developed in the previous section contain terms (e.g. X_B) that represent the aerodynamic forces acting on the vehicle. In this section these terms are defined as functions of the vehicle's state.

Although more sophisticated techniques are available (see the discussion in Reference 2), for the purposes of rigid body six degree-of-freedom simulation, it is usually quite adequate to use a quasisteady aerodynamic model based on Bryan's aerodynamic derivative technique.

No attempt has been made here to generalize this model so that it applies equally well to all types of air vehicles. Rather, its form has been simplified so that it is suitable for use only with free-flight, ballistic, rocket-boosted vehicles.

The resulting model expressed as aerodynamic force and moment components in F_B is summarized below. No attempt is made to rationalize this model other than to state that its use has resulted in predicted trajectories that are in good agreement with measured flight characteristics (see, e.g., Reference 1) of CRV7/BATS and ROBOT-9 vehicles.

The aerodynamic forces are specified by

$$X_{A_B} = -C_D q_D S \qquad (2.4,1a)$$

$$Y_{A_B} = C_{y\beta}\beta q_D S + q_D S (C_{y\alpha fin} + C_{y_{ofin}})_{pseudo}$$
 (2.4,1b)

$$Z_{A_B} = C_{z\alpha} \alpha q_D S + q_D S (C_{z\alpha fin} + C_{z_{ofin}})_{pseudo}$$
 (2.4,1c)

The various quantities in these equations are defined in the notation. It is important to note that q_E is the dynamic pressure given by

$$q_D = \frac{1}{2} \varrho V^2 \tag{2.4,2}$$

where ϱ is the air density and V is the airspeed given by

$$V = (u_B^2 + v_B^2 + w_B^2)^{1/2}$$
 (2.4,3)

Here (u_B, v_B, w_B) are the airspeed vector components expressed in F_B , and in the presence of nonzero wind conditions will not be the same as $(u_{B_E}, v_{B_E}, w_{B_E})$. Rather, they will be related to the wind velocity vector components in F_B [$(u_{B_g}, v_{B_g}, w_{B_g})$] through the relationships

$$u_B = u_{B_E} - u_{B_R} (2.4,4a)$$

$$v_B = v_{B_E} - v_{B_g} (2.4,4b)$$

$$W_B = W_{B_E} - W_{B_o} \tag{2.4,4c}$$

The latter may be obtained by considering the fundamental vector relationship

$$\underline{\mathbf{V}}_{\mathcal{E}} = \underline{\mathbf{V}} + \underline{\mathbf{W}} \tag{2.4,5}$$

i.e. the ground velocity vector $\underline{\underline{V}}_{\varepsilon}$ equals the airspeed vector $\underline{\underline{V}}$ plus the wind velocity vector $\underline{\underline{W}}$.

 C_D is the nondimensional drag coefficient, $C_{\nu\beta}$ is the aerodynamic derivative relating y-force due to sideslip angle β , $C_{z\alpha}$ is the aerodynamic derivative relating z-force due to angle of attack α , and S is a reference area that is usually the fuselage cross-sectional area for ballistic vehicles. C_D , $C_{\nu\beta}$, and $C_{z\alpha}$ may, in general, be Mach number and Reynolds number dependent, although here they are considered to be only Mach number dependent. α and β are given with respect to the x-axis of F_B , and from geometric considerations may be shown to be (see Figure 2)

$$\alpha = \arctan(w_B/u_B) \tag{2.4,6a}$$

$$\beta = \arctan(v_B/V_{xx}) \tag{2.4,6b}$$

where

$$V_{xx} = (u_B^2 + w_B^2)^{1/2} (2.4,7)$$

Differential equations may be obtained for α and β by differentiating (2.4,6a) and (2.4,6b) with respect to time, with the results

$$\dot{\alpha} = [\dot{\mathbf{w}}_B/\mathbf{V} - \dot{\mathbf{u}}_B \dot{\mathbf{w}}_B/\mathbf{V}^2] \cos^2 \alpha \qquad (2.4.8a)$$

and

$$\dot{\beta} = [V_{xz}\dot{v}_B - v_B(u_B\dot{u}_B + w_B\dot{w}_B)V_{xz}]/(V_{xz}^2\cos\beta) \qquad (2.4.8b)$$

The last terms on the right hand sides of (2.4,1b) and (2.4,1c) are pseudo fin terms and are included to permit modeling of aerodynamic asymmetries (e.g. due to production tolerances). They do not include any of the effects produced by the vehicle's nominal fin configuration. The latter have already been included in C_p , $C_{y\beta}$ and $C_{z\alpha}$.

The pseudo fin terms are defined as follows:

$$C_{\nu_{ofin}} = -C_{L_{ofin}} \delta_{fin} \sin \phi_{fin}$$
 (2.4,9a)

$$C_{x_{ofin}} = -C_{L_{ofin}} \delta_{fin} \cos \phi_{fin}$$
 (2.4,9b)

$$C_{yafin} = -e_{fin}C_{L_{\alpha}fin}\sin\phi_{fin} \qquad (2.4,10a)$$

$$C_{z_{\alpha fin}} = -e_{fin}C_{L_{\alpha fin}}\cos\phi_{fin} \qquad (2.4,10b)$$

$$\mathbf{w}_{fin} = \mathbf{w}_{B} \cos \phi_{fin} + \mathbf{v}_{B} \sin \phi_{fin} \tag{2.4,11}$$

$$\alpha_{fin} = \arctan(w_{fin}/u_B) \tag{2.4,12}$$

Here $C_{L_{\alpha}fin}$ is the lift slope of the fin, δ_{fin} is the cant angle of the fin, ϕ_{fin} is the angular cylindrical coordinate of the fin, $e_{f.n}$ is a body-fin interference factor and α_{fin} is the angle of attack of the fin. The fin geometry coordinate system is summarized in Figure 4.

The aerodynamic moment expressions are defined similarly to the aerodynamic force expressions, as follows:

$$L_{A_B} = C_{l_p} p_B q_D Sb^2 / (2V) + C_{l \delta fin} \delta_{fin} q_D Sb - Y_{A_{R_{nom}}} Z_{cg} - Z_{A_{R_{nom}}} Y_{cg}$$
(2.4,13a)

$$M_{A_B} = Z_{A_{Bnom}}(x_{ac} - x_{cg}) + C_{m_q} q_B S q_D b^2 / (2V) + X_{A_B} Z_{cg} + q_D S (C_{z_{\alpha fin}} \alpha_{fin} + C_{z_{\alpha fin}})_{pseudo} (x_{ac_{fin}} - x_{cg})$$
(2.4,13b)

$$N_{A_B} = -Y_{A_{Bnom}}(X_{ac} - X_{cg}) + C_{n_r} r_B S q_D b^2 / (2V)$$

$$+ X_{A_B} y_{cg} - q_D S (C_{yqfin} + C_{yqfin})_{pseudo} (X_{pcfin} - X_{cg})$$
(2.4,13c)

In these expressions $Y_{A_{Bnom}}$ and $Z_{A_{Bnom}}$ are given by (2.4,1b) and (2.4,1c) without the pseudo fin contributions, (x_{cg}, y_{cg}, z_{cg}) are the coordinates of the vehicle centre-of-mass in the vehicle structural reference frame F_R (see Figure 2), (x_{ac}, y_{ac}, z_{ac}) are the coordinates of the vehicle aerodynamic centre in F_R , $(x_{zc_{fin}}, y_{ac_{fin}}, z_{ac_{fin}})$ are the coordinates of the aerodynamic centre of the pseudo fin in F_R , b is the reference length (usually the fuselage diameter for ballistic vehicles) and the aerodynamic derivatives C_{l_p} , C_{ldfin} , C_{m_q} , $C_{z\alpha fin}$, C_{n_r} , $C_{y\alpha fin}$ are defined in the notation list.

The aerodynamic forces and moments may also be written as components in F_B by making an identical set of substitutions into (2.4,1) and (2.4,13) as used in converting the equations of motion written in F_B to those written in F_B' . It is important to note that certain simplifications result because of the underlying assumptions used in developing

the $F_{B'}$ equations, i.e. no pseudo fin terms may be included and y'_{B} and z'_{B} axes characteristics are identical.

For the sake of brevity, the aerodynamic force and moment expressions in F'_B will not be given here.

2.5 Mass and Moments of Inertia Models

The equations of motion have been written so that variations in the vehicle's mass and inertia characteristics (due to rocket motor propellant burn) are permitted. Component methods are used to compute the total vehicle mass and moments of inertia. The components considered are the vehicle airframe, the vehicle payload, the vehicle rocket motors less propellant and the rocket motors' propellant. Of these components, only the propellant characteristics are considered to be variable with time. Finally, the assumption has been made that the payload, the rocket motors, and the propellant are point masses.

Under these conditions the expressions for the vehicle mass and inertia characteristics are summarized below. These equations are given for the reference frame F_B . Position coordinates are with respect to the vehicle structural reference frame F_B . The subscripts used to reference the different components are as follows:

- 1) 'em' airframe (empty)
- 2) 'PL' payload
- 3) 'Me' rocket motors less propellant
- 4) 'PR' rocket motor propellant

 N_M is the total number of locket motors. Other variables used are defined precisely in the notation.

The expressions for the mass and inertia characteristics are given by $(\triangle x_{\xi} = x_{\xi} - x_{cg}, \Delta y_{\xi} - y_{\xi} - y_{cg}, \text{ and so forth})$

$$m = m_{em} + m_{PL} + \sum_{i=1}^{N_M} [(m_{M_e})_i + (m_{PR})_i]$$
 (2.5,1)

$$I_{xx}^{B} = I_{xx_{em}}^{B} + m_{em}(\Delta y_{em}^{2} + \Delta z_{em}^{2}) + m_{PL}(\Delta y_{PL}^{2} + \Delta z_{PL}^{2})$$

$$+ \sum_{i=1}^{N_{M}} \{ (m_{M_{e}})_{i} [(\Delta y_{M_{e}})_{i}^{2} + (\Delta z_{M_{e}})_{i}^{2}] + (m_{PR})_{i} [(\Delta y_{PR})_{i}^{2} + (\Delta z_{PR})_{i}^{2}] \}$$
(2.5,2a)

$$I_{yy}^{B} = I_{yy_{em}}^{B} + m_{em}(\Delta x_{em}^{2} + \Delta z_{em}^{2}) + m_{PL}(\Delta x_{PL}^{2} + \Delta z_{PL}^{2})$$

$$+ \sum_{i=1}^{N_{M}} \{(m_{M_{e}})_{i}[(\Delta x_{M_{e}})_{i}^{2} + (\Delta z_{M_{e}})_{i}^{2}] + (m_{PR})_{i}[(\Delta x_{PR})_{i}^{2} + (\Delta z_{PR})_{i}^{2}]\}$$
(2.5,2b)

$$I_{zz}^{B} = I_{zz_{em}}^{B} + m_{em}(\Delta x_{em}^{2} + \Delta y_{em}^{2}) + m_{PL}(\Delta x_{PL}^{2} + \Delta y_{PL}^{2}) + \sum_{i=1}^{IJ_{M}} \{ (m_{M_{e}})_{i} [(\Delta x_{M_{e}})_{i}^{2} + (\Delta y_{M_{e}})_{i}^{2}] + (m_{PR})_{i} [(\Delta x_{PR})_{i}^{2} + (\Delta y_{PR})_{i}^{2}] \}$$
(2.5,2c)

$$I_{x_{\nu}}^{B} = - m_{em} \Delta x_{em} \Delta y_{em} - m_{PL} \Delta x_{PL} \Delta y_{PL}$$

$$- \sum_{i=1}^{N_{M}} \{ (m_{M_{e}})_{i} (\Delta x_{M_{e}})_{i} (\Delta y_{M_{e}})_{i} + (m_{PR})_{i} (\Delta x_{PR})_{i} (\Delta y_{PR})_{i} \}$$
(2.5,2d)

$$I_{xx}^{B} = - m_{em} \Delta x_{em} \Delta z_{em} - m_{PL} \Delta x_{PL} \Delta z_{PL}$$

$$- \sum_{i=1}^{N_{M}} \{ (m_{M_{e}})_{i} (\Delta x_{M_{e}})_{i} (\Delta z_{M_{e}})_{i} + (m_{PR})_{i} (\Delta x_{PR})_{i} (\Delta z_{PR})_{i} \}$$
(2.5,2e)

$$I_{yz}^{B} = - m_{em} \Delta y_{em} \Delta z_{em} - m_{PL} \Delta y_{PL} \Delta z_{PL}$$

$$- \sum_{i=1}^{N_{M}} \{ (m_{M_{e}})_{i} (\Delta y_{M_{e}})_{i} (\Delta z_{M_{e}})_{i} + (m_{PR})_{i} (\Delta y_{PR})_{i} (\Delta z_{PR})_{i} \}$$
(2.5,2f)

$$X_{cg} = \{ m_{em} X_{em} + m_{PL} X_{PL} + \sum_{i=1}^{N_M} [(m_{me})_i (X_{me})_i + (m_{PR})_i (X_{PR})_i] \} / m$$
 (2.5,3a)

$$y_{cg} = \{m_{em}y_{em} + m_{PL}y_{PL} + \sum_{i=1}^{N_M} [(m_{me})_i(y_{me})_i + (m_{PR})_i(y_{PR})_i]\}/m \qquad (2.5,3b)$$

$$z_{c_R} = \{m_{e_m}y_{e_m} + m_{PL}y_{PL} + \sum_{i=1}^{N_M} [(m_{me})_i(z_{me})_i + (m_{PR})_i(z_{PR})_i]\}/m \qquad (2.5,3c)$$

Because it has been assumed that the only mass changes are due to propellant burn, in these expressions the only time variable quantities will be $(m_{PR})_i$, $(x_{PR_i}, y_{PR_i}, z_{PR_i})$. If we further assume that the propellant burns in such a way that the centre-of-mass of the propellant of a given rocket motor does not change significantly (e.g. as would be the case in rocket motors that are **not** end burners), then $(x_{PR_i}, y_{PR_i}, z_{PR_i})$ are not time variable and only $(\dot{m}_{PR})_i$ need be considered. The latter is related to the specific impulse of the rocket motor through the relationship (Reference 5)

$$(\dot{m}_{PR})_i = T_i(t)/(I_{sp_i}g)$$
 (2.5,4)

where g is the acceleration due to gravity, $T_i(t)$ is the thrust of the i-th rocket motor as a function of time t, and I_{sp_i} is the specific impulse of the i-th motor.

Under these assumptions and with (2.5,4), \dot{m} , the moment of inertia time derivatives, and $(\dot{x}_{cg}, \dot{y}_{cg}, \dot{z}_{cg})$ may be readily computed. For the sake of brevity, an exhaustive set of equations will not be given. Typically we have

$$\dot{m} = \sum_{i=1}^{N_M} \left[-T_i(t)/(I_{sp_i}g) \right]$$
 (2.5,5)

$$\dot{I}_{xx}^{B} = -2m_{em}(\Delta y_{em}\dot{y}_{cg} + \Delta z_{em}\dot{z}_{cg}) - 2m_{PL}(\Delta y\dot{y}_{cg} + \Delta z_{PL}\dot{z}_{cg})
+ \sum_{i=1}^{N_{M}} \{-2(m_{M_{e}})_{i}[(\Delta y_{M_{e}})_{i}\dot{y}_{cg} + (\Delta z_{M_{e}})_{i}\dot{z}_{cg}] + (\dot{m}_{PR})_{i}
[(\Delta y_{PR})_{i}^{2} + (\Delta z_{PR})_{i}^{2}] - 2(m_{PR})_{i}[(\Delta y_{PR})_{i}\dot{y}_{cg} + (\Delta z_{PR})_{i}\dot{z}_{cg}]\}$$
(2.5,6a)

$$\dot{\mathbf{I}}_{xy}^{B} = \mathbf{m}_{em}(\dot{\mathbf{x}}_{cg} \Delta \mathbf{y}_{em} + \Delta \mathbf{x}_{em} \dot{\mathbf{y}}_{cg}) + \mathbf{m}_{PL}(\dot{\mathbf{x}}_{cg} \Delta \mathbf{y}_{PL} + \dot{\mathbf{y}}_{cg} \Delta \mathbf{X}_{PL})
- \sum_{i=1}^{N_{M}} \left\{ -(\mathbf{m}_{M_{e}})_{i} \left[\dot{\mathbf{x}}_{cg} (\Delta \mathbf{y}_{M_{c}})_{i} + \dot{\mathbf{y}}_{cg} (\Delta \mathbf{X}_{M_{e}})_{i} \right] + (\dot{\mathbf{m}}_{PR})_{i} \right. (2.5,6b)
- (\Delta \mathbf{X}_{PR})_{i} (\Delta \mathbf{y}_{PR})_{i} - (\mathbf{m}_{PR})_{i} \left[\dot{\mathbf{x}}_{cg} (\Delta \mathbf{y}_{PR})_{i} + \dot{\mathbf{y}}_{cg} (\Delta \mathbf{X}_{PR})_{i} \right] \right\}$$

$$\dot{x}_{cg} = \left\{ \sum_{i=1}^{N_M} \left[-T_i(t)/(I_{sp_i})g\right](x_{PR})_i - x_{cg}\dot{m} \right\} m^{-1}$$
(2.5,7)

The mass and moment of inertia characteristic formulations expressed in the nonrotating reference frame F_{B} ' follow from the F_{B} expressions by incorporating the simplifying assumptions used for writing the F_{E} ' equations of motion (see Section 2.3). In particular, we have

$$I_{xy}^B = I_{xz}^B = I_{yz}^B = 0 (2.5,8)$$

and

$$I_{yy}^{B} = I_{zz}^{B} (2.5,9)$$

Equations (2.5,8) and (2.5,9) are just the result of the mass and inertia axisymmetry assumption used in developing the $F_{B'}$ equations of motion.

For the sake of brevity, the expressions for the $F_{B'}$ mass and inertia characteristics will not be given explicitly.

2.6 Thrust Characteristics

In Section 2.3 the equations of motion were written in the reference F_B with the thrust forces and moments written generally as $(X_{T_B}, Y_{T_B}, Z_{T_B})$ and $(L_{T_B}, M_{T_B}, N_{T_B})$ respectively. In this section these terms are examined in more detail.

The force terms $(X_{T_B}, Y_{T_B}, Z_{T_B})$ depend on the time domain thrust characteristics and the physical location and orientation of the rocket motors. This data must be known a priori to the simulation and is provided as input data to the computer program in the form of the thrust versus time look-up tables. The transformations used are summarized in Appendix 2.

The thrust moments require a somewhat more detailed examination. They are considered to consist of two components:

- 1) A moment due to the location and orientation of the thrust vector relative to F_B (see Figure 4),
- 2) A moment induced due to fixed vanes or nozzle grooves onto which the exhaust jet impinges.

Thus we have

$$L_{T_R} = (L_{T_B})_{c_g} + (L_{T_B})_{n_z}$$
 (2.6,1a)

$$M_{T_B} = (M_{T_B})_{i_R} + (M_{T_B})_{n_Z}$$
 (2.6,1b)

$$N_{T_B} = (N_{T_B})_{c_o} + (N_{T_B})_{nz}$$
 (2.6,1c)

These characteristics are recket motor specific. It is assumed that such data is available for the rocket motors used in the simulation. It then follows that once the orientation and location of the rocket motor thrust vectors relative to the vehicle are specified, enough information is available to determine $(L_{r_B}, M_{r_B}, N_{r_B})$ as given by (2.6,1a) to (2.6,1c) (a detailed treatment is given in Appendix 2 of Volume 2).

An assumption that has tacitly been made in this description of the thrust effects is that Coriolis forces and moments on the vehicle generated by the rocket motor exhaust are negligible. This need not always be the case, particularly for the moments, if the exhaust mass flow rate \dot{m} and the exhaust velocity vector relative to the vehicle V_R are large. However, for vehicles in the class of CRV7/BATS and ROBOT-9 using short burn duration 70 mm (2.75 inch) rocket motors, these effects are negligible and will not be considered further in this report.

2.7 Vehicle Kinematic Restrictions While on Launcher

The presence of the launcher during the initial portion of the flight places a number of kinematic constraints on the vehicle's motion. This section considers these constraints for a rail launcher such as was used for CRV7/BATS and ROBOT-9 (see Reference 1).

The basic geometrical quantities are defined in Figure 5. The equations of motion while the vehicle is on the rail are presented for the following assumptions:

- 1) The vehicle is mechanically constrained from tipping burnwards or forwards by the guide T-bolt until the T-bolt clears the launch rail (i.e. the vehicle is initially constrained to move along the x-axis of reference frame F_L). In the case of CRV7/BATS and ROBOT-9 there is also the launcher cage constraining the vehicle for part of its travel on the launch rail (see Reference 1).
- 2) The quantity s_G represents the distance the vehicle must move in the λ -direction of F_B before the guide T-boh clears the launch rail. The

bolt is assumed to be back far enough on the vehicle so that no significant tip-off may occur after it is clear of the rail and prior to the whole vehicle coming clear.

3) The vehicle may not move backwards on the launcher rail (i.e. \dot{u}_B is never less than zero).

Under these assumptions it follows that if the distance that the vehicle centre-ofmass has travelled (s) is less than or equal to s_G , then the vehicle is physically constrained to move only in the launch rail direction, i.e. for $s \le s_G$ we have

$$\dot{p}_B = \dot{q}_B = \dot{r}_B = \dot{v}_B = \dot{w}_B = 0$$
 (2.7,1a)

$$\dot{\mathbf{u}}_B \geq 0 \tag{2.7,1b}$$

$$\mathbf{v}_{B} = \mathbf{v}_{B}(0) \tag{2.7.2a}$$

$$\mathbf{w}_B = \mathbf{w}_B(0) \tag{2.7.2b}$$

$$p_B = p_B(0)$$
 (2.7,2c)

$$q_B = q_B(0) (2.7,2d)$$

$$r_B = r_B(0) (2.7,2e)$$

The nonzero conditions (2.7,2a) to (2.7,2e) allow for a nonstationary launcher, i.e. as would be the case for a launch from a ship in linear and angular motion.

The quantity s is defined precisely as

$$s \stackrel{\triangle}{=} \sqrt{(x_I - x_{I_o})^2 + (y_I - y_{I_o})^2 + z_I^2}$$
 (2.7,3)

where (x_I, y_I, z_I) are the vehicle centre-of-mass coordinates in F_I and $(x_{I_0}, y_{I_0}, 0)$ are the centre-of-mass coordinates when the vehicle is at rest on the launcher prior to first stage ignition.

For $s > s_G$, the governing equations are the unconstrained equations of motion developed in Section 2.3.

2.8 Wind Model

The aerodynamic model presented in Section 2.4 includes the wind velocities in F_B (F_B') as (u_{B_g} , v_{B_g} , w_{B_g}) [(u_{B_g} , v_{B_g}' , w_{B_g}')], and has tacitly assumed that there is no variation of the wind velocity from one point to another. This is equivalent to assuming that the wind induced aerodynamic loads are determined by its velocity rate of change acting at the centre-of-mass of the vehicle, an assumption referred to as the uniform-gust approximation (References 2 and 3). This approximation is equivalent to assuming that the wind velocity spectral content significantly affecting the vehicle response is at wavelengths that are greater than the significant vehicle dimensions (Reference 3). This assumption is reasonable when considering the rigid body dynamic response of flight vehicles, particularly for smaller vehicles such as ROBOT-9 and CRV7/BATS.

The wind velocity vector components relative to F_I are most conveniently expressed as components in F_I , i.e. (W_1, W_2, W_3) . These components may then be related to the wind velocity components in F_B with the rotation matrix \underline{L}_{BI} as given by (2.2,2a), i.e.

$$(u_{B_g}, v_{B_g}, w_{B_g})^T = \underline{L}_{BI}(W_1, W_2, W_3)^T$$
 (2.8,1)

The simulation package has provisions for inputting (W₁, W₂, W₃) as functions of altitude. This allows modeling of wind velocity atmospheric boundary layer effects, vehicle encounters with jetstream regions, and so forth. As well, since meteorological winds aloft data is usually given as a function of altitude, simulation of measured wind conditions is facilitated.

2.9 Atmospheric Conditions

Since the ROBOT-9 and CRV7/BATS vehicles have the capability to achieve altitudes well above 9000 m (30,000 ft), an atmospheric model is required that takes into account variations in density (ϱ), temperature (T_A), pressure (p_A), and the speed of sound (a) as a function of altitude above sea level (h_{ASL}).

The models used are based on the U.S. standard atmosphere (1962), as is common practice in aeronautical engineering, and are valid within the troposphere, i.e. for $h_{ASL} \le 11,100 \text{ m}$ (36,000 ft) (see Reference 6). They are given by

$$T_A = 288.15 - 0.0065 h_{ASL} R$$
 (2.9,1)

$$P_A = 101300.(T_A/288.15)^{5 \cdot 255} (2.9,2)$$

$$a = 20.0463 \sqrt{T_A} \tag{2.9,3}$$

$$\varrho = 0.00348454 \, p_A / T_A \tag{2.9,4}$$

where

$$R = r_E/(h_{ASL} + r_E) (2.9,5)$$

 T_A is in degrees Kelvin, p_A is in Pascals, a is in meters per second, ϱ is in kilograms per meter cubed, h_{ASL} is in meters, and r_E is the Earth's radius to the sea level datum,

$$r_e = 6.3567658 \times 10^6 \text{ m} (2.0855531 \times 10^7 \text{ ft}).$$

From the factor R it is also convenient to compute the variation of the acceleration due to gravity as a function of altitude, i.e.

$$g = g_o R^2 (2.9,6)$$

where $g_o = 9.80667 \text{ m/s}^2 (32.1741 \text{ f/s}^2)$.

Provision has been made in the simulation package to vary the temperature and pressure (and thus the density) from the standard values by allowing altitude dependent per cent deviations from standard conditions.

2.10 Aspect Angle Equations

For target and flight test applications, it is frequently necessary that the vehicle's aspect azimuth (ξ_A) and elevation (ξ_E) angles be known with respect to an observer at F_T (see Figure 1). This section presents equations for ξ_A and ξ_E in terms of the location and orientation of F_T relative to that of F_I .

It is assumed that the x-y planes of F_I and F_T are parallel, i.e. that F_I may be rotated to F_T through a rotation ψ_T about the z-axis of F_I .

Let the vector position of F_T relative to F_I be \underline{R}_{TI} , and that of the vehicle centre-of-mass relative to the origin of F_I be \underline{R} (see Figure 6). It follows that the vector position of the vehicle relative to F_T is given by

$$\underline{\mathbf{R}}_{T} = \underline{\mathbf{R}} - \underline{\mathbf{R}}_{TI} \tag{2.10,1}$$

or in matrix notation

$$\underline{\mathbf{R}}_{T}^{T} = \underline{\mathbf{L}}_{TI} \underline{\mathbf{R}}^{I} - \underline{\mathbf{L}}_{TI} \underline{\mathbf{R}}_{TI}^{I} \tag{2.10,2}$$

where \underline{L}_{TI} is the rotation matrix rotating vector components in F_I to components in F_T , and is given by (2.2,4a).

Equations (2.10,2) may be written in scalar form as

$$x_T = (x_I - x_{TI})\cos\psi_T + (y_I - y_{TI})\sin\psi_T$$
 (2.10,3a)

$$y_T = -(x_I - x_{TI})\sin\psi_T + (y_I - y_{TI})\cos\psi_T$$
 (2.10,3b)

$$z_T = z_I - z_{TI}$$
 (2.10,3c)

From the definition of ξ_A and ξ_E in Figure 1, it follows that

$$\dot{\xi}_A = \arctan(y_T/x_T) \tag{2.10,4a}$$

$$\xi_E = \arctan\left(-z_T/x_T\right) \tag{2.10,4b}$$

3. BALSIM SOFTWARE DESCRIPTION — GENERAL

The dynamic model described in the previous chapter has been implemented in the BALSIM simulation package. All coding was carried out using IBM FORTRAN for the H-extended compiler. The package has been debugged and tested on the IBM 3033 computer, and has been used to predict the dynamic characteristics of the CRV7/BATS and ROBOT-9 vehicles.

The software is currently being installed on VAX11/780 and Honeywell DPS-8/70C computers.

The software consists of a MAIN program plus nine subroutines making up approximately 885 FORTRAN source statements. There are no subroutines or functions, other than these, that are not available in standard FORTRAN on-line libraries.

The software userbook is given in Appendix 1 of Volume 2 with a source language listing of the package.

3.1 Software Capabilities

The dynamic model implemented with the BALSIM package has already been discussed in detail in the previous chapter. Its limitations will not be considered further here.

The package was developed with the objective of providing a convenient basis for inputting the characteristics of multistaged rocket vehicles and predicting their dynamic rigid body characteristics. By appropriately modifying the input data set, it provides for

- 1) nominal and off-nominal vehicle mass, inertia and thrust characteristics,
- 2) different motor types,
- 3) Mach number dependent aerodynamic characteristics,
- 4) structural production tolerances,
- 5) system failures (e.g. stage and fin failures),
- 6) moving launchers,
- 7) user specified initial conditions,
- 8) user specified payload characteristics,
- 9) tabular output in either metric or English units, and
- 10) multiple case runs.

As well, with minor software modification, response calculations may be stored on disk for subsequent use with other software (e.g. plotting software).

3.2 Software Cimitations

Lift its current form the software is not intended for use in the following types of simulations:

- 1) Ballistic rocket vehicles with control surfaces.
- 2) Winged flight vehicles.
- 3) Nonrigid vehicles,
- 4) Simulations where the inertial flat Earth approximations are invalid.
- 5) Vehicles where the staging process involves physically releasing rocket motor stages.

Limitations (1) and (5) may be removed with relatively minor alterations to the dynamic model of Chapter 2 with corresponding changes to the software.

3.3 Numerical Integration Algorithm

The numerical integration algorithm used to solve the system of ordinary differential equations describing the vehicle's dynamics is a fixed step-size, fourth order Runge-Kutta method (see Reference 7). Provision has been made for the specification of two step sizes, one for use during rocket motor burns, and the other for use during coasting flight. The latter technique was found to considerably reduce CPU time in certain simulations.

3.4 Software Testing and Execution Times

The BALSIM package has been used extensively to predict the performance and dynamic characteristics of the CRV7/BATS and ROBOT-9 vehicles. These predictions have been used to define the nominal, dispersion and safe y-envelope characteristics (see References 8 and 9) of these vehicles.

Flight test data obtained early in the development of CRV7/BATS (see Reference 1) was used to update and validate the aerodynamic model that has been employed. More recent comparisons with flight test data have also proved to be satisfactory (also see Reference 1).

BALSIM predictions have also been evaluated for consistency by comparing results obtained using the equations of motion written in F_B with those written in $F_{B'}$, with satisfactory results.

The execution CPU time of the package will depend on the computer used, on the step sizes chosen, and on the duration of the flight time simulated. For the IBM 3033 computer, the following CPU execution times were observed for simulations of the ROBOT 9 vehicle using a 0.05 second integration step size, and a 1.0 second tabulated output increment:

- 1) For 8 cases averaging 108 simulated flight seconds per case, the execution CPU time was 343.4 seconds, yielding a 0.4 seconds CPU execution time per simulated flight second ratio.
- 2) The compilation CPU time with the H-extended compiler was 25 seconds.
- 3) The linkage editor CPU time was 1.6 seconds.

This completes the general description of the BALSIM software package. Detailed user related data is given in Volume 2.

4. SUMMARY

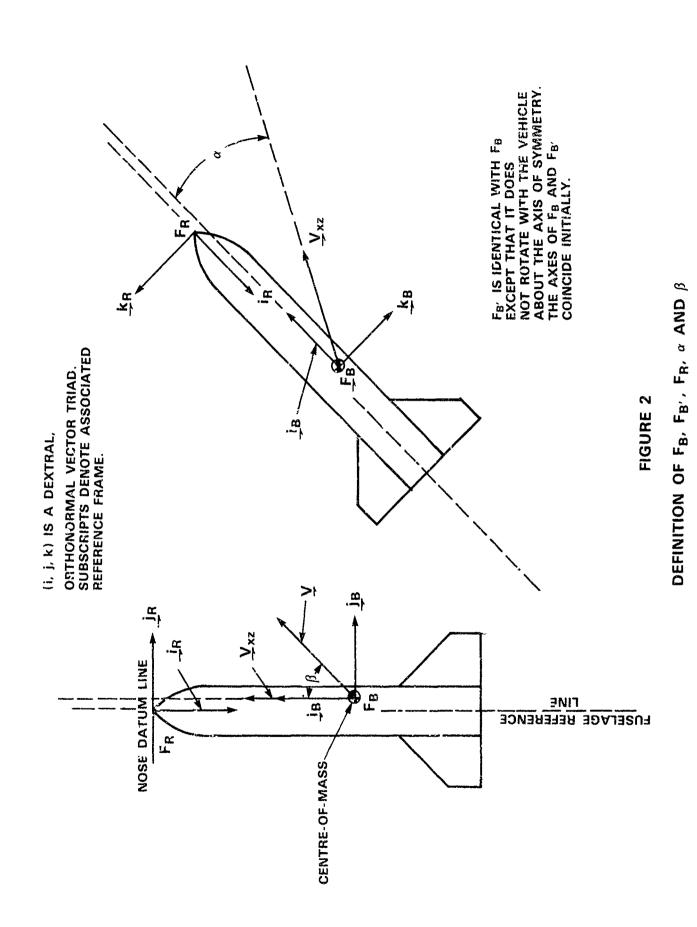
Six degree-of-freedom, rigid body equations of motion suitable for modeling the dynamic characteristics of multistaged, free-flight, ballistic rockets have been rigorously developed, and have been implemented in a FORTRAN software package called BALSIM. This package allows for modeling of vehicle thrust and structural asymmetries, time-varying mass and inertia characteristics, variable wind conditions, nonstandard atmospheric conditions, stage failures, and different rocket motor types.

The EALSIM package has been successfully used to predict the performance and dynamic characteristics of the CRV7/BATS and ROBOT-9 vehicles both with and without moving launchers. It will be adapted for use with the VAX11/780 and Honeywell DPS-8 computers in the near future.

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DEFINITION OF FL, FL, FT



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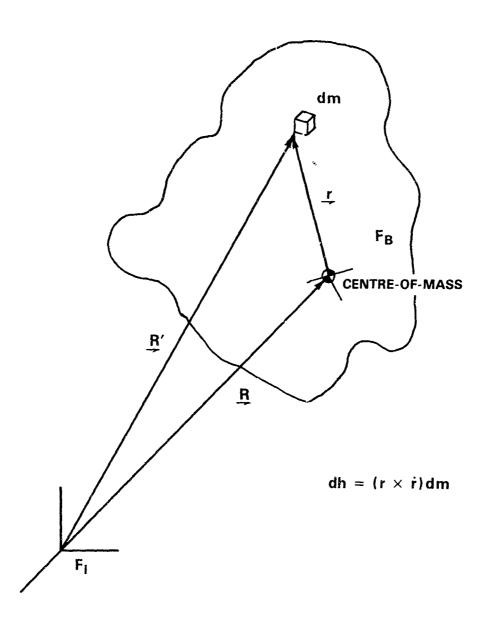


FIGURE 3

ANGULAR MOMENTUM CONTRIBUTION OF
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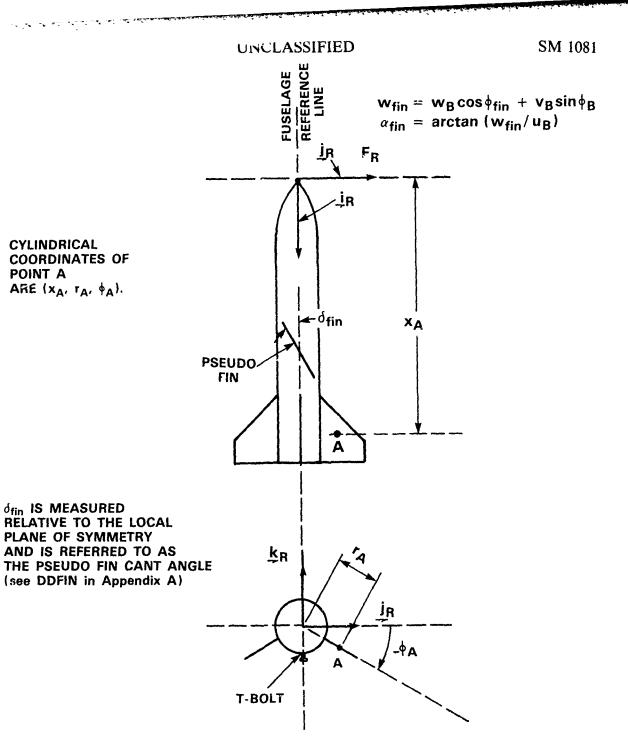


FIGURE 4
VEHICLE CYLINDRICAL COORDINATE SYSTEM

s = distance vehicle has moved along launcher rail relative to rest position

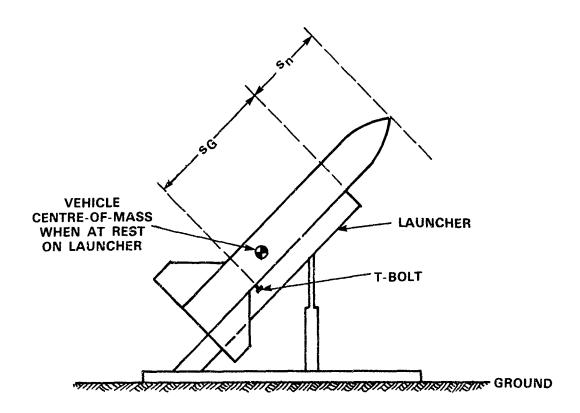
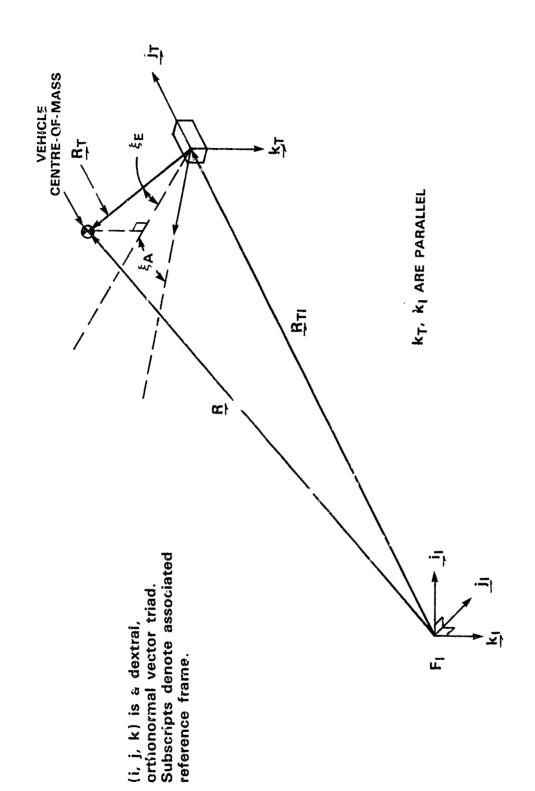


FIGURE 5

VEHICLE KINEMATIC CONSTRAINTS WHILE ON
LAUNCHER GEOMETRY



ASPECT ANGLE GEOMETRY

FIGURE 6

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Six degree-of-freedom, rigid body equations of motion are described suitable for modeling the dynamic characteristics of multistaged, free-flight, ballistic rockets such as the DPES developed aerial targets CRV7/BATS and ROEOT-9. These equations of motion form the core of a FOPTRAN simulation software package called BALSIM. This package allows for modeling of vehicle thrust and structural asymmetries, time-varying wass and inertia characteristics, variable wind conditions, nonstandard atmospheric conditions, stage failures, and different rocket motor types. The BALSIM package has been written in IBM FORTRAN IV and has been tested on the IBM 3033 computer with the H-extended compiler. It is currently being adapted for use with the VAX11/780 and Honeywell DPS-8/70C computers. (II)

KEY WORDS

Aerial Targets CRV7/BATS Flight Dynamics ROBOT-9 Stability and Control Six degree-of-freedom simulation

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